Ultracompact dwarfs in the Perseus Cluster: UCD formation via tidal stripping

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ABSTRACT

We present the results of a Keck/DEIMOS survey of Ultracompact dwarfs (UCDs) in the Perseus Cluster core. We confirm cluster membership for 14 UCDs, with radial velocities $\sim 5300$ km s$^{-1}$. Two of these confirmed that Perseus UCDs have extremely blue colours $(B - R < 0.6$ mag), reside in star-forming filaments surrounding NGC 1275, and have likely formed as massive star clusters in the last $\sim 100$ Myr. We also measure a central velocity dispersion of a third, UCD13 ($\sigma_0 = 38 \pm 8$ km s$^{-1}$), the most extended UCD in our sample. We determine it to have radius $R_e = 85 \pm 1$ pc, a dynamical mass of $(2.3 \pm 0.8) \times 10^8$ M$_\odot$, and a metallicity $[Z/H] = -0.52^{+0.33}_{-0.20}$ dex. UCD13 and the cluster’s central galaxy, NGC 1275, have a projected separation of 30 kpc and a radial velocity difference of $\sim 20$ km s$^{-1}$. Based on its size, red colour, internal velocity dispersion, dynamical mass, metallicity, and proximity to NGC 1275, we argue that UCD13 is likely the remnant nucleus of a tidally stripped dwarf elliptical (dE), with its progenitor dE having $M_B \approx -16$ mag and mass $\sim 10^9$ M$_\odot$.

Key words: galaxies: clusters: individual: Perseus Cluster – galaxies: dwarf – galaxies: star clusters: general.

1 INTRODUCTION

Ultracompact dwarfs (UCDs) were first identified over a decade ago (Hilker et al. 1999; Drinkwater et al. 2000), but whether they should be considered as galaxies or massive star clusters is still uncertain (e.g. Hassel et al. 2005; Hilker 2009; Mieske et al. 2009; Forbes & Kroupa 2011; Willman & Strader 2012). Brodie et al. (2011) adopted a working definition for UCDs to have half-light sizes of $10 < R_h < 100$ pc and $M_V < -8$ mag (with no upper limit on luminosity).

The two main formation channels discussed in the literature are variations of merged star clusters (Kroupa 1998; Fellhauer & Kroupa 2002; Brüns et al. 2011) or stripped dwarf galaxies (Bassino, Muzzio & Rabolli 1994; Bekki, Couch & Drinkwater 2001a). In this sense, their current size and luminosity are thought to be the result of a ‘bottom-up’ growth from lower masses, or a ‘top-down’ evolution from higher masses. Support for the idea of both formation channels for UCDs comes from the results of Norris & Kannappan (2011).

They found many UCDs to be consistent with high-mass globular clusters (GCs). However, they also reported one UCD around NGC 4546 that was in a counter-rotating orbit with respect to the host galaxy and its GC system. Furthermore, they calculated that the UCD had a very short dynamical friction time-scale. This strongly suggested that it had been recently captured and was not part of the galaxy’s original GC system.

Evidence for UCDs forming as massive star clusters in the nearby Universe was presented by Penny, Forbes & Conselice (2012), who found young, blue candidate UCDs with $(B - R)_0 < 0.6$ mag in the Perseus Cluster associated with gas filaments around the cluster’s central galaxy NGC 1275. NGC 1275 therefore provides a location in which to study the formation of massive star clusters and proto-UCDs in the nearby Universe. If these massive blue star clusters are able to remain bound against the cluster tidal potential, they will be seen as typical red UCDs once their stellar populations fade in a time-scale of $\sim 500$ Myr.

Using a state-of-the-art compilation of pressure-supported systems with old stellar populations, Brodie et al. (2011) found that UCDs had a tight colour–magnitude relation that was different from normal GCs. This, and kinematic differences, argued against the
merged star cluster scenario. However, they also found that UCDs do not have a well-defined size–luminosity distribution as might have been expected in the stripping scenario (Pfeffer & Baumgardt 2013). A further expectation of the stripping scenario is that the remnants come from a range of progenitor masses. Hence, we would expect a continuity of objects between the most massive UCDs and compact ellipticals (cEs), the latter of which are thought to be the tidally stripped remnant of a normal galaxy (Faber 1973). Such intermediate-mass objects are too massive to be explained as merged star clusters.

In the compilation of Brodie et al. (2011), their fig. 8 shows very few objects with confirmed sizes and luminosities intermediate between those of UCDs and cEs. One such object is M59cO highlighted by Chilingarian & Mamon (2008) as a ‘missing link’ between UCDs and cEs. This object, with a confirmed velocity placing it in the Virgo Cluster, has $M_B = -12.34$ mag ($M_V \sim -13.0$ mag). A dual Sérsic fit to the surface brightness profile gave a ‘nucleus’ of half-light radius $R_e = 13$ pc with an extended component of $R_e = 50$ pc. The object was also found to be old and relatively metal rich. Another UCD in the Virgo Cluster, M60-UCD1 (Strader et al. 2013) with $M_B = -13.24$ mag and a dynamical mass $M_{dyn} = (2.0 \pm 0.3) \times 10^8 M_\odot$ is the densest galaxy identified to date and occupies the region of size–magnitude space between the most massive GCs and the faintest cE galaxies.

In this work, we present the results of a Keck-DEIMOS follow-up study of UCD candidates in the nearby Perseus Cluster ($D = 71$ Mpc) identified in Penny et al. (2012) out to 250 kpc from the cluster centre. The brightest cluster galaxy (BCG) is the unusual NGC 1275, which is actively accreting/disrupting satellite galaxies at the present time, making Perseus an ideal environment in which to study the formation of UCDs as disrupted dwarf ellipticals (dEs) versus massive star clusters. The focus of this paper is in which to study the formation of UCDs as disrupted dwarf ellipticals versus massive star clusters. The origin of UCDs in clusters: Are they massive star clusters or the cores of tidally stripped galaxies? In particular, we examine on the origin of the largest Perseus UCD, UCD13, via its size, kinematics, cluster location, and stellar population.

This paper is organized as follows. In Section 2, we describe our observations and their reduction. In Section 3, we briefly discuss the properties of the confirmed Perseus Cluster members, including the size–luminosity distribution for UCDs. An analysis of UCD13 is presented in Section 4 with a tidal stripping formation scenario for UCD13 examined in Section 5. A star cluster origin for UCD13 is discussed in Section 6. We discuss our results in Section 7 and present our conclusions in Section 8.

2 DATA

2.1 HST ACS imaging

In Penny et al. (2012), from Hubble Space Telescope (HST) Advanced Camera for Surveys (ACS) WFC imaging, we identified 84 candidate UCDs in the Perseus Cluster core using two fields in the F475W ($B$) and F625W ($R$) bands, along with five fields in the F555W ($V$) and F814W ($I$) bands. These UCD candidates were selected based on their sizes alone ($10 < R_e < 150$ pc), with no colour or luminosity cuts placed on their selection. Instead, all candidate UCDs were identified to have $S/N > 40$ to allow for an accurate determination of their size via ISHAPE (Larsen 1999), along with smooth, round shapes with no evidence for sub-structure. All UCD candidates identified using this method were brighter than $M_B = -10.5$ mag. The reduction and analysis of the HST data are described in Penny et al. (2012). The sizes, magnitudes, and colours of all Perseus UCDs presented in this paper are taken from Penny et al. (2012), except for the size of UCD13 which we remeasure here. These UCD properties are shown in Table 1.

<table>
<thead>
<tr>
<th>Object</th>
<th>$R_e$ (pc)</th>
<th>$R$ (mag)</th>
<th>$(B-R)_0$ (mag)</th>
<th>$V$ (km s$^{-1}$)</th>
<th>Line used</th>
</tr>
</thead>
<tbody>
<tr>
<td>UCD1</td>
<td>11.7</td>
<td>22.81 ± 0.02</td>
<td>1.18 ± 0.05</td>
<td>5717 ± 10.9</td>
<td>CaT</td>
</tr>
<tr>
<td>UCD2</td>
<td>26.9</td>
<td>22.90 ± 0.03</td>
<td>1.44 ± 0.06</td>
<td>5687 ± 25.2</td>
<td>CaT</td>
</tr>
<tr>
<td>UCD7</td>
<td>12.5</td>
<td>22.04 ± 0.02</td>
<td>1.41 ± 0.04</td>
<td>5272 ± 15.3</td>
<td>CaT</td>
</tr>
<tr>
<td>UCD9</td>
<td>14.4–24.3</td>
<td>22.85 ± 0.02</td>
<td>1.26 ± 0.05</td>
<td>9955 ± 12.1</td>
<td>CaT</td>
</tr>
<tr>
<td>UCD13*</td>
<td>85.0</td>
<td>20.97 ± 0.01</td>
<td>1.37 ± 0.03</td>
<td>5292 ± 13.7</td>
<td>CaT</td>
</tr>
<tr>
<td>UCD17</td>
<td>18.4</td>
<td>22.97 ± 0.03</td>
<td>0.28 ± 0.04</td>
<td>5286 ± 17.3</td>
<td>Hα (emission)</td>
</tr>
<tr>
<td>UCD19</td>
<td>13.1</td>
<td>21.88 ± 0.02</td>
<td>1.17 ± 0.03</td>
<td>5680 ± 11.1</td>
<td>CaT</td>
</tr>
<tr>
<td>UCD21</td>
<td>31.2</td>
<td>22.23 ± 0.02</td>
<td>1.10 ± 0.04</td>
<td>5668 ± 10.9</td>
<td>CaT</td>
</tr>
<tr>
<td>UCD23</td>
<td>11.5</td>
<td>22.36 ± 0.02</td>
<td>1.14 ± 0.05</td>
<td>5693 ± 14.5</td>
<td>CaT</td>
</tr>
<tr>
<td>UCD25</td>
<td>11.5</td>
<td>21.66 ± 0.01</td>
<td>1.54 ± 0.03</td>
<td>4750 ± 14.5</td>
<td>CaT</td>
</tr>
<tr>
<td>UCD29</td>
<td>22.4</td>
<td>22.80 ± 0.02</td>
<td>0.54 ± 0.04</td>
<td>5250 ± 16.5</td>
<td>Hα (emission)</td>
</tr>
<tr>
<td>UCD40</td>
<td>13.9</td>
<td>22.12 ± 0.02</td>
<td>1.47 ± 0.04</td>
<td>5632 ± 10.2</td>
<td>CaT</td>
</tr>
<tr>
<td>UCD60</td>
<td>49.0</td>
<td>21.43 ± 0.02</td>
<td>0.96 ± 0.04</td>
<td>3735 ± 15.8</td>
<td>CaT</td>
</tr>
<tr>
<td>UCD65</td>
<td>13.6</td>
<td>22.05 ± 0.02</td>
<td>1.05 ± 0.05</td>
<td>7681 ± 15.3</td>
<td>CaT</td>
</tr>
</tbody>
</table>

*The effective radius for UCD13 is the result of a single Sérsic fit to its light distribution. All other sizes are the result of King profile fits.

Table 1. Perseus Cluster UCDs with spectroscopically confirmed cluster membership. Lower and upper limits are presented for the size of UCD9 (see the explanatory text in Section 2.2). Apparent magnitudes, colours, and effective radii are taken from Penny et al. (2012). The effective radii presented here are determined from King profiles. The radial velocities are obtained from KECK-DEIMOS spectroscopy. The spectral feature used for the radial velocity determination is also shown. Magnitudes and colours for UCDs 1 through 40 are in $B$ and $R$, and those for UCD60 and UCD65 are in $I$ and $V$ reflecting the available HST imaging.
Of the 14 UCD candidates with good spectra, we are able to confirm 13 of them as belonging to the Perseus Cluster (mean velocity of 5366 km s$^{-1}$), with one (UCD9) likely located just outside of the main cluster. This implies a background contamination rate in our HST-selected UCD candidates (Penny et al. 2012) of $\leq$7 per cent. The candidate UCD9 has a measured velocity of $9555 \pm 40.5$ km s$^{-1}$. Relative to the velocity dispersion of the elliptical galaxies in the cluster (i.e. 1324 km s$^{-1}$), it is 3.5$\sigma$ from the mean. However, relative to the dwarf galaxies (1818 km s$^{-1}$), it is only 2.5$\sigma$ (see Penny & Conselice 2008). Thus, we conclude that UCD9 is likely associated with the large-scale filamentary structure/infall regions surrounding the Perseus Cluster, and include it in our sample of Perseus UCDs. We take the size of $R_e = 14.4$ pc for UCD9 presented in Penny et al. (2012) as a lower limit on UCD9’s effective radius. We calculate an upper limit for the UCDs’ size if it is instead background to Perseus using the online cosmology calculator of Wright (2006), assuming $H_0 = 71$ km s$^{-1}$ Mpc$^{-1}$. At $z = 0.033$ (the redshift of UCD9), the angular scale is 590 pc arcsec$^{-1}$ versus 350 pc arcsec$^{-1}$ for Perseus. The upper limit on the size of UCD9 is $R_e = 24.3$ pc, compared to the size $R_e = 14.4$ pc presented in Penny et al. (2012) and calculated assuming the UCD is located in the Perseus Cluster. Both the upper and lower limit on UCD9’s size are presented in Table 1.

The effective radii of the Perseus UCDs were calculated in Penny et al. (2012) by converting the full width at half-maximum (FWHM) of their King profile (King 1962) to an effective radius using the conversion $R_e = 1.48 \times$ FWHM. While King profiles are generally a good fit to the light profiles of GCs and small UCDs, they are less robust for more extended UCDs. Of the UCDs in the Penny et al. (2012) sample, one object, UCD13 is extended enough that we are able to measure its $R_e$ directly with a Sérsic fit to its light distribution (Section 4.1).

In Fig. 2, we show the size–luminosity distribution for all 15 confirmed Perseus UCDs, along with confirmed old-aged (>5 Gyr) pressure-supported systems from the compilation of Brodie et al. (2011), which was updated by Forbes et al. (2013). The data presented in Fig. 2 are taken from the following sources: Barmby et al. (2007), Brassere et al. (2011), Brodie & Larsen (2002), Chilingarian et al. (2009), Cockcroft et al. (2011), Da Costa et al. (2009), Fadely et al. (2009), and many others.
et al. (2011), Forbes et al. (2013), Foster et al. (2011), Geha et al. (2003), Harris (2010), Hasçegan et al. (2005), Hasçegan (2007), Hau et al. (2009), Hwang et al. (2011), Huxor et al. (2011), Martin et al. (2008), Mieske et al. (2007), Misgeld & Hilker (2011), Mouhcine et al. (2010), Norris & Kannappan (2011), Richtler et al. (2005), Rejkuba et al. (2007), Romanowsky et al. (2009), Smith Castelli et al. (2008), Strader et al. (2011), and van den Bergh & Mackey (2004). For a full description of the data compilation used here, see Misgeld et al. (2011), Brodie et al. (2011), Forbes et al. (2013) and references therein.

Where necessary, values of $M_V$ were converted to $M_B$ using the colour $(B - V) = 0.96$ mag for a typical elliptical galaxy taken from Fukugita, Shimasaku & Ichikawa (1995). This colour transformation is comparable to the $(B - V)$ colours of UCDs in the Fornax Cluster (Karick, Drinkwater & Gregg 2003) and Coma Cluster (Chiboucas et al. 2011), which exhibit colours ranging between $(B - V) = 0.65$ and 1.1 mag.

Figure 3. $F625W$ ($R$) band HST ACS image of UCD13. A low surface brightness stellar envelope is clearly seen around the UCD.

In general, the sizes of the confirmed Perseus UCDs agree with those of UCDs in other studies. Furthermore, the sizes of the smaller UCDs with $R_e < 20$ pc and $M_B \sim -10.5$ mag are in good agreement with those of the Côté et al. (2006) sample of dE nuclei. This suggests that a fraction of these objects could be the nuclei of stripped dEs, that have had their fragile stellar envelopes removed by the tidal potential of the Perseus Cluster, as well as during encounters with other cluster members. These smaller UCDs also overlap in size–luminosity space with the most extended GCs, so multiple origins are possible for these UCDs. The 14 confirmed UCDs cover a range of colours. Thus, as well as confirming the status of several old stellar population UCDs, we have confirmed that two of the very blue UCDs (UCD17 and UCD29) with $(B - R) < 0.6$ mag are associated with the Perseus Cluster and must be young massive star clusters.

The largest Perseus UCD, UCD13, is approaching the region of the size–luminosity diagram occupied by cEs, and the object exhibits a low surface brightness stellar envelope (Fig. 3). An effective radius of 85 pc makes it one of the most extended UCDs identified to date, more extended than a typical GC by a factor of 30. At $M_B = -11.92$ mag, it is furthermore the brightest UCD with confirmed Perseus Cluster membership. Given its large size and brightness, UCD13 is an interesting object in which to examine the origin of extended UCDs; its origin as a massive, extended star cluster is more difficult to explain than the smaller Perseus UCDs with effective radii $< 20$ pc. We now go on to examine UCD13 in more detail by considering its light profile, kinematics, and stellar population.

4 UCD13: ANALYSIS

4.1 Surface brightness profile

Surface photometry for UCD13 is fitted to HST ACS $F625W$ ($\sim R$) band imaging using the IRAF task ISOPHOTE ELLIPSE. IRAF ELLIPSE fits elliptical isophotes to a galaxy image at pre-defined semi-major axis values. The isophote centres, ellipticities, and position angles are allowed to vary during the fit, with the result of this fit shown in Fig. 4. The UCD is remarkably round at all radii, with ellipticity $\epsilon < 0.14$ out to a radius of 280 pc. The position angle profile for the
UCD does not exhibit any sharp twists that would indicate internal sub-structure such as spiral arms or a bar, but instead gradually varies out to the last reliable isophote at a semi-major axis of 280 pc. The UCD’s F435W – F625W (B – R) colour profile is flat within the error bars, suggesting that the UCD does not exhibit a metallicity gradient.

We model UCD13’s light distribution with GALFIT (Peng et al. 2002), using the form of the Sérsic profile as defined in Caon, Capaccioli & D’Onofrio (1993):

\[
\mu(R) = \mu_e + \frac{2.5 b_n}{\ln(10)} \left( \frac{R}{R_e} \right)^{-n} - 1,
\]

(1)

where \( R \) is the semi-major axis, \( R_e \) is the effective radius, and \( n \) is the Sérsic index. The constant \( b_n \) is defined in terms of the Sérsic index as \( b_n = 1.9992 n - 0.3271 \) for the range 0.5 < \( n < 10 \).

The analytic profile for the UCD is convolved with the HST ACS point spread function (PSF) prior to the modelling of the UCD. The PSF is constructed using TINY TIM (Krist 1993) at the position of UCD13 on the ACS detector to take into account the distortions that are present in the original image. GALFIT then constructs a model of the UCD that best fits the data by minimizing the residuals between the model and the original input image.

We first fit UCD13 with a Sérsic model. A single Sérsic profile fit has \( R_e = 85 \pm 1.1 \) pc, \( n = 2.99 \pm 0.03 \), and \( \mu_{R_625W} = 21.59 \pm 0.02 \) mag arcsec\(^{-2} \). However, the single Sérsic profile is a poor fit to the data, with large residuals between the measured and model surface brightness (\( \Delta \mu_{R_625W} > 0.1 \) mag arcsec\(^{-2} \), especially at large radii. UCDs are commonly found to be best fitted with two-component models (e.g. Chilingarian et al. 2011), so we investigate if this is the case for UCD13.

For a two-component fit, the inner profile has \( R_e = 22 \pm 0.2 \) pc, \( n = 0.67 \pm 0.02 \), \( \varepsilon = 0.90 \pm 0.01 \), and \( \mu_{R_625W} = 19.5 \pm 0.01 \) mag arcsec\(^{-2} \). The outer component has \( R_e = 144 \pm 2.1 \) pc, \( n = 1.28 \pm 0.01 \), \( \varepsilon = 0.90 \pm 0.01 \), and \( \mu_{R_625W} = 22.67 \pm 0.02 \) mag arcsec\(^{-2} \). The results of this fit, along with the residuals, are shown in the left-hand plot of Fig. 4. UCD13 is well fitted by a two-component Sérsic model, with typical residuals \( \Delta \mu_{R_625W} < 0.05 \) between the surface brightness profile of UCD13 and the PSF-convolved model of UCD13 constructed using GALFIT.

### 4.2 Internal kinematics

To help establish if this object is the stripped remnant of a much larger system, we measure its internal velocity dispersion using the penalized pixel fitting routine (Cappellari & Emsellem 2004). The instrumental velocity dispersion is confirmed using two stellar spectra that were taken as serendipitous observations during the DEIMOS observations (these stars fell into the slits of our target objects).

We measure UCD13 to have an internal velocity dispersion \( \sigma = 35 \pm 8 \) km s\(^{-1} \), determined from the CaT 1 and CaT 2 lines, and confirmed using the H\( \alpha \) line. The CaT line was affected by a strong sky line, so was not used in the determination of \( \sigma \). We then correct this global internal velocity dispersion to a central velocity dispersion \( \sigma_0 \) (the velocity dispersions within one \( R_e \)). Using an effective radius \( R_e = 85 \) pc as found for a single profile fit, we obtain \( \sigma_0 = 38 \pm 9 \) km s\(^{-1} \).

We calculate the dynamical mass of the UCD using the virial theorem. To calculate the total mass of the object, we assume an exponential surface brightness profile (\( n = 1 \)) as a compromise between the inner and outer profiles, and a half-light radius \( R_e = 85 \) pc. Using the virial theorem with a correction for the Sérsic profile shape provided in Cappellari & Emsellem (2004), this gives us a dynamical mass for the object of \((2.3 \pm 0.8) \times 10^8 M_\odot \).

We estimate a stellar mass-to-light ratio (\( M/L \)) using the stellar population synthesis models of Bruzual & Charlot (2003). We assume a Salpeter IMF, an old stellar population with an age of 10 Gyr, and a metallicity between \( Z = 0.004 \) and 0.008 estimated from the metallicity presented in Section 4.3. Based on this, we find a B-band \( M/L \) of \( 5 \pm 1 \). The absolute B-band magnitude of UCD13 \( M_B = -11.92 \) is converted to a solar luminosity and then multiplied by the stellar \( M/L \), giving a stellar mass for the object of \((4.4 \pm 0.9) \times 10^9 M_\odot \). This gives \( M_{dyn}/M_* \) of 5.2. However, the IMF for UCDs may be better described by a Kroupa IMF (e.g. Chilingarian et al. 2011). Using a Kroupa IMF with [Z/H] = -0.33 and an old stellar population (Maraston 2005), the B-band stellar \( M/L \) ratio will decrease to 3.66, resulting in \( M_{dyn}/M_* = 7.2 \). Both results for \( M_{dyn}/M_* \) are inconsistent with a highly dark-matter-dominated object.

### 4.3 Metallicity

We use the method of Usher et al. (2012) to measure the metallicity of UCD13 from the strength of the near-infrared CaT spectral feature. Sensitive to metallicity and insensitive to ages older than 3 Gyr (Vazdekis et al. 2003), the strength of CaT is independent of \( \alpha \)-element enhancement and horizontal branch morphology (Brodie et al. 2012). However, the strength of the CaT is dependent on the IMF as the CaT is weaker in a stellar population with a bottom-heavy IMF, such as those found in massive early-type galaxies by van Dokkum & Conroy (2010), than in a population with a Kroupa IMF and the same metallicity (Vazdekis et al. 2003).

To remove the effect of strong sky-line residuals on the CaT measurement, we mask the sky-line regions and fit a linear combination of stellar templates to the observed spectrum. We then measure the strength of CaT on the fitted spectrum and use a Monte Carlo resampling technique to estimate a confidence interval on the measured line strength. Lastly, we use the single stellar population models of Vazdekis et al. (2003) to convert the line strength into a metallicity. See Usher et al. (2012) for further details of the metallicity measurement process.

We measure a metallicity of \([Z/H] = -0.52^{+0.33}_{-0.29} \) dex from the strength of the CaT. Using the GC colour–metallicity relation of Usher et al. (2012), we derive a colour-based metallicity of \([Z/H] = -0.13 \pm 0.23 \) dex, consistent with the CaT-based metallicity. The metallicity and other properties of UCD13 measured in this work are summarized in Table 2, along with three comparison objects from the literature (M59e, Chilingarian & Mamon 2008; VUCD7, Evstigneeva et al. 2007; and M60-UCD1, Strader et al. 2013).

### 5 UCD13: TIDALLY STRIPPED GALAXY OR STAR CLUSTER?

Could UCD13 be the tidal remnant of a once much larger galaxy? UCD13 has a velocity \( v = 5292 \pm 13.7 \) km s\(^{-1} \), compared to \( v = 5276 \) km s\(^{-1} \) for NGC 1275 (Petrosian et al. 2007). At an angular separation of 1.5 arcmin, UCD13 and NGC 1275 have a minimum physical separation of 30 kpc. Such conditions are highly favourable for galaxy–galaxy interactions and tidal stripping. NGC 1275 furthermore exhibits a shell structure consistent with the BCG actively disrupting and accreting lower mass galaxies today (e.g. Penny et al. 2012), and UCD13’s progenitor likely
contribute to these shells. It is therefore highly likely that the two objects are interacting given their small radial velocity separation of <20 km s\(^{-1}\) and close physical proximity, so we examine UCD13 for features that might indicate that it is a UCD formed via the tidal stripping of a low-mass galaxy.

5.1 The velocity dispersion–luminosity relation

When a galaxy is tidally stripped, its effective radius and luminosity will be reduced, but its central velocity dispersion and central metallicity will remain largely unchanged from those of its progenitor (Bender, Burstein & Faber 1992). As the tidal remnant relaxes after stripping, its effective radius will increase slightly, causing its central velocity dispersion to decrease; however, this effect will be small. The simulations of Chilingarian et al. (2009) show that for a disc galaxy, the stellar disc of the galaxy will be heavily stripped, and its central velocity \(\sigma_0\) will decrease by <5 km s\(^{-1}\). This decrease in velocity dispersion is within our error bar of ±8 km s\(^{-1}\), and we can assume that the progenitor and remnant galaxies will have the same value of \(\sigma_0\).

For the tidal stripping of a nucleated dE with a fragile stellar envelope (i.e. most of its mass is located in its central regions), we hypothesize that this decrease in \(\sigma_0\) will be even smaller. In addition, as shown in Fig. 5, at a given magnitude dEs exhibit a scatter in their measured internal velocity dispersions, with \(\sigma_0\) ranging from ~20 to 70 km s\(^{-1}\). Any decrease in \(\sigma_0\) resulting from the tidal transformation of a dE to a UCD, given our measured value of \(\sigma_0 = 38\) km s\(^{-1}\), will fall within this range. We can therefore use our measurement of UCD13’s central velocity dispersion to characterize the nature of its progenitor.

Old, pressure-supported galaxies follow a well-defined relation between \(\sigma\) and luminosity: the Faber–Jackson relation (Faber & Jackson 1976). However, the slope of this relation is observed to change in the dE regime, from \(L \propto \sigma^4\) for giant ellipticals to \(L \propto \sigma^2\) for dEs (Matković & Guzmán 2005). This allows us to more easily separate the kinematics of giants and dEs. We can therefore use the luminosity–sigma relation to estimate the mass of the progenitor of UCD13 by examining its location on the velocity dispersion–luminosity relation (Fig. 5), with a high value of \(\sigma_0\) likely indicative of a tidal stripping origin.

UCD13’s central velocity dispersion is comparable to the internal velocity dispersion of a dE galaxy with \(M_B = -16\) mag (\(\sigma_0 = 38\) km s\(^{-1}\), Fig. 5). We can rule out UCD13 as currently being a dSph with a high-surface-brightness core, as its velocity dispersion is lower than that of dSph galaxies of comparable magnitude. In the Local Group, Fornax, Sculptor, and Leo I have luminosities comparable to UCD13 (−12 mag < \(M_V < -10\) mag), but central velocity dispersions \(\sigma < 12\) km s\(^{-1}\) (Gilmore et al. 2007; Walker et al. 2007). The UCD’s \(\sigma_0\) is too low for it to be the remnant of a low-mass elliptical galaxy or bulge. We confirm the luminosity of UCD13’s progenitor by utilizing the relations of Matković & Guzmán (2005). Using the relations presented in their figs 4 and 5, \(\sigma_0 = 38\) km s\(^{-1}\) corresponds to a galaxy with \(M_B = -17.5\) mag and \(B - R = 1.35\) mag. UCD13 has \(B - R_B = 1.37\) mag, in strong agreement with this predicted progenitor colour. Therefore, based on colour and internal velocity dispersion, the progenitor of UCD13 is likely a dE galaxy with \(M_B = -16\) mag.
We should note than an internal velocity dispersion of \( \sim 40 \text{ km s}^{-1} \) is not inconsistent with the velocity dispersion--magnitude relation for massive star clusters. However, the size of UCD13 \((R_e = 93 \text{ pc})\) is higher than that expected for an extended star cluster with a comparable \( \sigma_0 \). Despite having a higher velocity dispersion of \( \sigma = 45 \pm 5 \text{ km s}^{-1} \), the massive luminous star cluster W3 has a size \( R_e = 17.5 \pm 1.8 \text{ pc} \) (Maraston et al. 2004), one-fifth the size of UCD13. For a comparable central velocity dispersion, massive star clusters are more compact than UCD13. A massive star cluster origin for UCD13 is examined in detail in Section 6.

5.2 Light distribution

The surface brightness distributions of dEs are generally well fitted by an exponential \( n \approx 1 \) Sérsic profile, while the light distributions of giant ellipticals closely follow a de Vaucouleurs profile with a Sérsic index \( n \approx 4 \). Although the luminosities of cEs frequently overlap those of dEs \((M_B > -18)\), the cEs are easily separated from dEs due to their concentrated light distributions from a single Sérsic profile fit, with \( n > 2 \). With \( n = 2.99 \) from a single profile Sérsic fit, UCD13 has a concentrated light distribution resembling that of a typical cE galaxy.

To establish if the object’s large size and luminosity could be the result of the tidal stripping of a dE, in Fig. 2, we include for comparison a sample of dE nuclei from Côte et al. (2006). UCD13 is more extended than a dE nucleus of comparable magnitude by a factor of 10. Evstigneeva et al. (2008) find a comparable result for UCDs in the Fornax and Virgo Clusters, with UCDs having effective radii larger than 20 pc, twice the size of those of early-type dE nuclei at the same luminosity.

However, during a tidal stripping event, the nucleus may be able to retain or re-accrete some of its envelope of tidally stripped stars. If a stripped dE is able to retain some of its stellar envelope in addition to its nucleus, the resulting UCD will be more extended than a dE nucleus alone. Pfeffer & Baumgardt (2013) show that UCDs with sizes \( R_e > 10 \text{ pc} \) can be formed from the tidal stripping of a nucleated dE if the dE has a small pericentric approach \(<10 \text{ kpc}\) from the perturbing galaxy. The dE must be on a highly elliptical orbit and must be on its first or second approach for such a transformation to occur. In their simulations, this scenario produces extended UCDs with sizes \( \sim 100 \text{ pc} \). The resulting UCDs consist not only of the remnant dE nucleus, but also some of the dE’s retained stellar envelope.

Given its close proximity to NGC 1275 \((\sim 30 \text{ kpc})\) and small separation in radial velocity \(<20 \text{ km s}^{-1}\), UCD13 would currently be near the apocentre of a radial orbit, else we are observing its orbit face-on (such an orbit could still be a radial). If UCD13 is indeed at an orbit apocentre \(\sim 30 \text{ kpc}\) after several pericentric approaches to NGC 1275, dynamical friction will have decayed its original highly elliptical orbit into a circular one, as predicted in the simulations of Chilingarian et al. (2011). Therefore, a dE stripping scenario is a feasible explanation for its origin.

5.3 The mass–metallicity relation

Assuming a dE rather than a massive cluster origin for UCD13, we proceed to examine the mass–metallicity relation for dwarf galaxies to estimate the mass of the UCD’s progenitor. The mass–metallicity relation for galaxies presented in fig. 6 of Mendel et al. (2009) shows that a galaxy with \([Z/H] = -0.52 \text{ dex} \) has a mass \(\sim 10^9 \text{ M}_\odot\). This mass is comparable to that of a dE with \(M_B \sim -16\), in agreement with the result using the Faber–Jackson relation presented in Section 5.1.

6 A MASSIVE STAR CLUSTER?

A galaxy-like central bulge with a stellar envelope as exhibited by UCD13 (Fig. 4) does not necessarily exclude a massive star cluster origin. The simulations of Fellhauer & Kroupa (2005) showed that UCDs with diffuse stellar envelopes similar to that observed for UCD13 can be formed when massive \( (>10^7 \text{ M}_\odot) \) star cluster complexes merge. Such objects will have circular isophotes in their inner regions, with their ellipticities increasing at large radii. Norris & Kannappan (2011) suggested \( M_V \sim -13 \text{ mag} \) as the upper magnitude limit for UCDs formed as giant GCs, with UCD13 approaching this magnitude limit assuming \((B - V) \sim 1\), a typical value for UCDs (e.g. Chiboucas et al. 2011). Multiple mergers between multiple massive star clusters would be required to reach this size, and the simulations of Fellhauer & Kroupa (2002, 2005) are able to form UCDs via this method.

UCDs formed out of a merged star cluster complex would have a core of merged star clusters, with an envelope of stripped stars lost during the merger of the star clusters – similar in morphology to UCD13. The intermediate-age star cluster W3 \((R_e = 17.5 \text{ pc})\) in the merger remnant NGC 7252 is likely a UCD-like object (Fellhauer & Kroupa 2005) recently formed via this method. However, its size is only one-quarter that of UCD13, so it is unclear if massive star clusters such as W3 are the same class of objects as the most extended UCDs. We therefore examine the sizes of massive star clusters surrounding NGC 1275 (e.g. Canning et al. 2010; Penny et al. 2012) to address this question.

NGC 1275 is surrounded by a complex system of filaments, a number of which host regions of active star formation. These star-forming regions are hypothesized as a region of UCD formation, with several proto-UCDs located in these filaments exhibiting \((B - V)_0\) colours \(<0.55\), implying ages \(<100 \text{ Myr} \). However, the sizes of these young UCD-like star clusters are far smaller than that found for UCD13, with the largest blue UCD, UCD29, having a size \(R_e = 22.4 \text{ pc} \) (Penny et al. 2012). Instead, objects such as UCD13 are more likely to form in lower density regions of the cluster. UCD13 and NGC 1275 are very close in both radial velocity and angular separation, implying UCD13 is located in the centre of Perseus, a region of the cluster that is not conducive to mergers between massive star clusters. A strong tidal field will prohibit merger events between massive star clusters. Canning et al. (2010) showed that a number of young, massive star clusters in the star-forming filaments surrounding NGC 1275 are being tidally stripped and will eventually fall into the galaxy’s nucleus, and thus NGC 1275’s filaments may not be an ideal location for forming more extended UCDs via star cluster mergers.

Mergers between multiple massive young, UCD-sized star clusters are required to produce an extended object such as UCD13. A star–cluster merger formation scenario for UCD13 therefore seems unlikely. However, other Perseus UCDs with sizes \( <20 \text{ pc} \) are likely a mix of massive GCs and galaxy nuclei.

7 DISCUSSION

The UCD population in Perseus may have multiple origins. UCDs with smaller effective radii are likely an extension of the GC system to higher magnitudes and sizes. However, the origin of UCDs more massive than this cannot be explained simply as massive GCs. Multiple mergers between massive GCs or star-forming complexes are
required to explain their origin. UCDs with galaxy-like properties and cEs likely represent a family of objects with a tidal stripping origin, with cEs similar to the Local Group object M32 being formed from the tidal stripping of low-luminosity spiral galaxies (Bekki et al. 2001b), while the most extended UCDs with sizes \(>40 \text{ pc}\) are probably stripped dEs. Massive UCDs would therefore simply be an extension of cE galaxies to lower masses and sizes into a regime where they overlap with the most massive GCs.

Based on its size and internal velocity dispersion, we have established that UCD13 is likely the tidal remnant of a larger progenitor. Similar objects to UCD13 are shown in Table 2, and include M59cO (Chilingarian & Mamon 2008) at \(M_B = -12.34\) mag with \(R_c = 32 \text{ pc}\), located at a distance of 9 kpc from the Virgo Cluster giant elliptical M59. This object has a blue core surrounded by a redder envelope, with an age \(9.3 \pm 1.4\) Gyr, and a metallicicity of \(Z = -0.03 \pm 0.04\) dex. Its internal velocity dispersion is \(28.3 \text{ km s}^{-1}\) (Norris et al., in preparation) and lies between the UCDs and cEs on the Fundamental Plane. Chilingarian & Mamon (2008) suggested that it is similar in metallicity to low-luminosity ellipticals and lenticulars in terms of age, metallicity, and effective surface brightness. UCD13 and M59cO have very similar properties, with M59cO exhibiting a two-component structure with a low surface brightness stellar envelope.

Chilingarian & Mamon (2008) suggested that M59cO is the result of tidal stripping by M59, with the UCD/cE transition object the remnant of a much larger object. Strader et al. (2013) reach a similar conclusion for the origin of the M60-UCD1, which is likely the densest dE with an \(M_B = -7.8\) mag located at a distance of \(1.37\) Mpc from the Virgo Cluster giant elliptical M60. This object has a blue core surrounded by a redder envelope, consistent with the simulations of Pfeffer & Baumgardt (2013). Its dynamical mass is \((2.3 \pm 0.8) \times 10^{6} \text{ M}_\odot\), making it one of the most massive UCDs identified to date. Its internal velocity dispersion, \(\sigma_v = 38 \pm 9 \text{ km s}^{-1}\), along with its metallicity \([Z/H] = -0.52^{+0.33}_{-0.20}\) dex and colour \((B - R)_b = 1.37\), are consistent with its origin as a dE with an \(M_B = -16\) and a mass \(\sim 10^{6} \text{ M}_\odot\). Given its size \((R_c = 85 \pm 1.1 \text{ pc})\) and location at the centre of Perseus, the UCD is inconsistent with being a massive star cluster.

8 CONCLUSIONS
We have examined the properties of Perseus UCDs with their cluster membership confirmed via KECK-DEIMOS spectroscopy. Perseus UCDs with sizes \(<20 \text{ pc}\) are consistent with being massive GCs or the nuclei of stripped dEs. Two blue UCDs with colours \((B - R) < 0.6\) are confirmed as members of the Perseus Cluster and must have formed within the last 100 Myr as massive star clusters in the star-forming filaments surrounding NGC 1275.

The largest UCD (UCD13) is likely the nuclear remnant of a tidally stripped dE that has retained a large fraction of its stellar envelope, consistent with the simulations of Pfeffer & Baumgardt (2013). Its dynamical mass is \((2.3 \pm 0.8) \times 10^{6} \text{ M}_\odot\), making it one of the most massive UCDs identified to date. Its internal velocity dispersion, \(\sigma_v = 38 \pm 9 \text{ km s}^{-1}\), along with its metallicity \([Z/H] = -0.52^{+0.33}_{-0.20}\) dex and colour \((B - R)_b = 1.37\), are consistent with its origin as a dE with an \(M_B = -16\) and a mass \(\sim 10^{6} \text{ M}_\odot\). Given its size \((R_c = 85 \pm 1.1 \text{ pc})\) and location at the centre of Perseus, the UCD is inconsistent with being a massive star cluster.

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