

# First Release of High-redshift Superluminous Supernovae from the Subaru HIgh-Z SUpernova CAmpaign (SHIZUCA). II. Spectroscopic Properties

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#### **Abstract**

We present Keck spectroscopic observations of three probable high-redshift superluminous supernovae (SLSNe) from the Subaru HIgh-Z sUpernova CAmpaign (SHIZUCA), confirming redshifts of 1.851, 1.965, and 2.399. The host galaxies were selected for transient monitoring from multiband photometric redshifts. The supernovae are detected during their rise, and the classically scheduled spectra are collected near maximum light. The rest-frame far-ultraviolet (~1000–2500 Å) spectra include a significant host galaxy flux contribution, and we compare our host-galaxy-subtracted spectra to UV-luminous SNe from the literature. While the signal-to-noise ratios of the spectra presented here are sufficient for redshift confirmation, supernova spectroscopic type confirmation remains inconclusive. The success of the first SHIZUCA Keck spectroscopic follow-up program demonstrates that campaigns such as SHIZUCA are capable of identifying high-redshift SLSNe with sufficient accuracy, speed, and depth for rapid, well-cadenced, and informative follow-up.

Key words: galaxies: high-redshift - methods: observational - supernovae: individual - ultraviolet: general

#### 1. Introduction

It is now known that some supernovae exceed an absolute magnitude of  $M \simeq -21$ , giving rise to a luminosity class of supernovae termed superluminous supernovae (SLSNe; Quimby et al. 2011; Gal-Yam 2012; Nicholl et al. 2015; Moriya et al. 2018). SLSNe are also more luminous in the far ultraviolet (FUV;  $\sim 1000-2500$  Å) relative to other supernovae (Quimby et al. 2011; Cooke et al. 2012; Howell et al. 2013; Yan et al. 2017b, 2018). SLSNe are exceedingly rare, occurring at  $\sim 0.001 \times$  the frequency of general core-collapse supernovae (Quimby et al. 2013; Prajs et al. 2017). But these transients promise to be powerful probes of the early universe as they are already visible in ground-based observations to redshifts of z=4 and greater (Cooke et al. 2012; Mould et al. 2017; Moriya et al. 2019).

Recent improvements to survey astronomy, such as large-format CCD mosaics, are enabling more wide-area surveys to operate, such as the Hyper Suprime-Cam Subaru Strategic Program (HSC-SSP; Aihara et al. 2018), the Dark Energy Survey (DES; Dark Energy Survey Collaboration et al. 2016), the Zwicky Transient Facility (Bellm & Kulkarni 2017), and Pan-STARRS (Kaiser et al. 2010). These surveys monitor the vast amounts of sky required to detect SLSNe with a reasonable

frequency. In addition, these surveys are broad enough to allow boutique surveys such as the Subaru HIgh-Z sUpernova CAmpaign (SHIZUCA; Moriya et al. 2019) and the Survey Using DECam for Superluminous Supernovae (SUDSS; Smith et al. 2016) to exercise more focused monitoring techniques and detect SLSNe with higher efficiency.

There are advantages to surveying for SLSNe at high redshift  $(z \ge 2)$ . Massive progenitors are invoked in all of the explosion mechanisms used to explain SLSNe (Gal-Yam 2012; Moriya et al. 2018), suggesting that the rate of SLSNe per unit volume varies with redshift as the cosmic star formation rate, rising and perhaps peaking at  $z \sim 2$  (Madau & Dickinson 2014). This is consistent with rates measured up to this redshift (Quimby et al. 2013; Prajs et al. 2017). Some if not all types of SLSNe show a preference for low-metallicity hosts (Lunnan et al. 2014; Leloudas et al. 2015; Angus et al. 2016; Schulze et al. 2018), suggesting that SLSNe may require low-metallicity progenitors and that the volumetric rate of SLSNe may continue to increase with redshift beyond the peak of star formation into epochs of universally low metallicity. The measured rate at  $z \sim 2-4$  is consistent with a continuing increase beyond  $z \sim 2$  (Cooke et al. 2012), though this measurement is based on two events.

Another advantage of looking to high redshift is the ability to sample the rest-frame FUV of SLSNe with optical facilities.

FUV spectra of low-redshift supernovae are rare because they can only be collected from space-based telescopes. Studies of such spectra have been done (Fransson et al. 2002, 2005, 2014; Panagia 2007; Bufano et al. 2009), including SLSNe (Yan et al. 2017b, 2018; Quimby et al. 2018). Due in part to their excessive UV luminosity and to model-predicted UV indicators of certain explosion mechanisms (Mazzali et al. 2016), there are programs focused on collecting FUV spectra of SLSNe in particular (Quimby 2014, 2016). Still, given the low rate of SLSNe, few such spectra have been collected to date. At high redshift, this critical UV information is pushed into the optical bands, enabling this wavelength region to be explored with ground-based observations. A number of FUV analyses of SLSNe have been accomplished in this way (Berger et al. 2012; Howell et al. 2013; Pan et al. 2017; Smith et al. 2018).

High éntendu surveys like SHIZUCA are already capable of identifying supernovae to z>4 (Mould et al. 2017; Moriya et al. 2019). Spectroscopic follow-up of high-redshift supernovae must be carried out on 8 m class telescopes due to their intrinsically faint apparent magnitudes (often  $m_r\gtrsim 24$  for  $z\gtrsim 2$ ).

Here we present spectra obtained with Keck of three probable high-redshift SLSNe. This paper summarizes the first spectroscopic follow-up program of SHIZUCA. It has been written in parallel with Moriya et al. (2019), the first photometric analysis of the SHIZUCA program, hereafter M19. We summarize the observations in Section 2. In Section 3 we present the spectra and their redshift measurements. In Section 4 we discuss our analysis of the observations, and we present our conclusions in Section 5. All calculations in this paper assume a  $\Lambda$ CDM cosmology with  $H_0 = 70 \, \mathrm{km \, s^{-1} \, Mpc^{-1}}, \, \Omega_M = 0.3, \, \mathrm{and} \, \Omega_\Lambda = 0.7.$  All magnitudes are AB and all wavelengths are quoted in rest-frame ångströms unless otherwise specified.

## 2. Observations

The COSMOS field (Scoville et al. 2007) of the HSC-SSP<sup>15</sup> supplies the photometry for this work and consists of a 1.8 deg<sup>2</sup> field of view Hyper Suprime-Cam pointing imaged in five filters (*grizy*). The first SHIZUCA season was active on the COSMOS field from 2016 November to 2017 May.

SHIZUCA uses COSMOS2015 photometric redshifts (Laigle et al. 2016) and MIZUKI redshifts (Tanaka 2015) to identify potential hosts of high-redshift transients from HSC-SSP photometry. The per-epoch depth of observation from the HSC-SSP using the 8.2 m Subaru telescope ( $m_i \lesssim 26.5$ ) enables SHIZUCA to monitor the fluxes of high-redshift sources in individual observing epochs. Any flux variations observed in these sources are then analyzed in the context of the photometric redshift of the host. Flux variations are identified as SLSN candidates using criteria such as nonrecurrence, duration, light curve shape, peak magnitude, and color evolution. Details of the SHIZUCA photometric analysis and relevant MIZUKI redshift probability distributions are presented and discussed in M19.

Follow-up spectroscopy of select SHIZUCA photometric high-redshift SLSN candidates was acquired on 2016 December 28 and 2017 March 22–23 using the Low Resolution Imaging Spectrometer (LRIS; Oke et al. 1995; Steidel et al. 2004) on the

Keck I telescope. These data were obtained using the 400/3400 grism on the blue arm and the 400/8500 grating on the red arm separated at  $\sim 5600$  Å using the D560 dichroic. The CCDs were read out using  $2\times 2$  binning with a spectral resolution of  $\sim 500$  km s<sup>-1</sup>. The FWHM seeing ranged from 0.9 to 1.11, and normal atmospheric extinction increased from light cirrus. During the March observations, the blue shutter of LRIS failed and was fixed to remain open, and the trapdoor was used in its place. The exposures were set at 1200 s on the blue side and 1179 s on the red side with the difference chosen to allow the CCDs to finish reading out at the same time.

For follow-up, we targeted SHIZUCA photometric candidates near maximum light projected to be  $m_r \lesssim 25.5$  during observation. This magnitude limit aims to achieve combined transient and host continuum signal-to-noise ratios (S/N) near rest-frame 1500 Å of S/N  $\sim 3-5$  pix<sup>-1</sup> for integrations of  $\sim$ 2 hr. SLSNe usually arise in star-forming galaxies such as Lyman break galaxies (LBGs), and these S/N are sufficient to reliably measure redshifts via identification of strong UV absorption features from the interstellar medium along with Ly $\alpha$  emission when present and the Ly $\alpha$  forest when visible (Steidel et al. 1998). However, the low S/N combined with systematics in the internal flats increase flux uncertainties at wavelengths shortward of 3600 Å in the observer frame.

On 2016 December 28, three SHIZUCA photometric candidates were eligible for follow-up, of which two were targeted in the available time, HSC16adga and HSC16aaqc. On 2017 March 22–23, another three SHIZUCA photometric candidates made up the primary target list: HSC17auzg, HSC17dbpf, and HSC17cywe. We followed up each of these targets with sufficient time remaining to follow up one backup target, HSC17davs, below our specified magnitude limit. A spectrum of such low S/N can still yield a measurable redshift and even supernova type information if it includes identifiable strong emission features.

The SHIZUCA photometric candidates are transient phenomena detected photometrically in or near host galaxies with high photometric redshifts (see M19). There are cases of photometric redshift confusion in which separate, discrete photometric redshift probabilities persist for a single galaxy. For example, the decrement in the rest-frame UV continuum of a  $z \sim 3$  LBG caused by the Ly $\alpha$  forest can be photometrically mimicked by the Balmer break in the rest-frame optical continuum of a  $z \sim 0.3$ galaxy (Steidel et al. 2003; Cooke et al. 2006). Because the number of targets available for follow-up in a narrow observation window is limited, candidates observed in hosts exhibiting photometric redshift confusion with a reasonable probability of being high redshift are not disqualified from our target lists. However, to minimize the loss of time observing low-redshift transients, we inspect the raw 2D spectra of each target as they are read out for any obvious emission features. Often one or two emission features are sufficient to differentiate between high- and low-redshift probabilities. If a target is determined to be low redshift during observation, the target is aborted and the next target is acquired.

From the six SHIZUCA photometric candidates followed up, three arise in z > 1.8 hosts, which can be reliably matched with an LBG reference spectrum (see Table 1 and Section 3). Two candidates, HSC16aaqc and HSC17cywe, were found to be low redshift through the positive association of emission features from the raw 2D spectra to commonly observed transitions in low-redshift galaxies (e.g., [O II]  $\lambda\lambda3727$ , H- $\beta$ , and [O III]  $\lambda5007$ ). The spectrum of one target, HSC17davs and host,

<sup>&</sup>lt;sup>15</sup> General information for the HSC-SSP, such as location, cadence, and data products, can be found at <a href="http://hsc.mtk.nao.ac.jp/ssp/">http://hsc.mtk.nao.ac.jp/ssp/</a> and the associated links

Table 1
SHIZUCA Follow-up Targets

SN-ID	Exp Time blue; red (s)	Photo-z <sup>a</sup>	Spec-z	λ Range <sup>b</sup> (Å)
2016 Dec HSC16adga HSC16aaqc	7200;7074 2400;2358	$2.26^{+0.25}_{-0.30} \\ 1.25^{+1.20}_{-0.03}$	2.399 ~0.36°	942–2765
2017 Mar HSC17auzg HSC17dbpf HSC17cywe HSC17davs	8400;8253 6000;5895 2400;2358 4800;4716	$1.65^{+0.06}_{-0.08} \\ 2.25^{+0.08}_{-0.53} \\ 0.45, 3.26^{d} \\ 3.34^{+1.27}_{-2.75}$	1.965 1.851 ~0.46° °	1080–3170 1123–3297 

#### Notes.

recorded no detectable signal or emissions. None of the spectra exhibited obvious signs of being other types of high-redshift transients, namely active galactic nuclei (AGNs) or tidal disruption events (TDEs).

Among the spectra we use as templates for comparison are those of the SLSN-II LSQ15abl (Quimby et al. 2015). The formal presentation and analysis of these spectra is part of a work in progress. They were collected with the *Hubble Space Telescope (HST)* using the Cosmic Origins Spectrograph (COS; Shull 2009) and the Space Telescope Imaging Spectrograph (STIS; Woodgate et al. 1997). The reduced spectra are from the Mikulski Archive for Space Telescopes (MAST) and have been adjusted for scale, noise, and redshift.

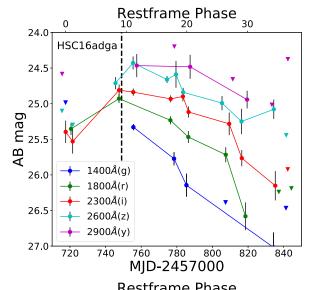
#### 3. Data Reduction and Analysis

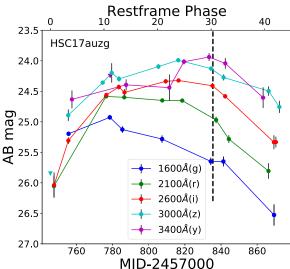
# 3.1. HSC16adga

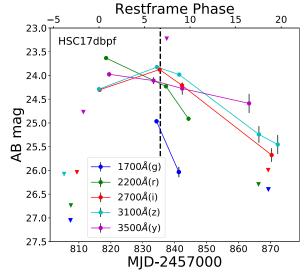
HSC16adga (SN-2016jhm) was first detected by SHIZUCA on MJD 57715 at coordinates (R.A., decl.) = (10:02:20.12, +02:48:43.4). The transient is detected in a star-forming host with a COSMOS photo-z of  $2.26^{+0.25}_{-0.30}$  and a MIZUKI photo-z peaking at 2.19. It is observed offset by  $0.^{\prime\prime}36\pm0.^{\prime\prime}104$  ( $\sim$ 3 kpc) from the host galaxy flux centroid, helping to rule out an AGN or a TDE. The light curve of HSC16adga is shown in Figure 1.

We collected the Keck LRIS spectrum of HSC16adga on MJD 57751.09. The positioning of the spectral slit is mapped in Figure 2. From the light curve in Figure 1, the spectrum appears to have been taken near peak in the rest-frame UV. The spectrum, composed of six exposures (7200 s blue, 7074 s red), was reduced using standard IRAF<sup>16</sup> procedures. The combined blue- and red-side 1D spectrum is shown in Figure 3.

The spectrum includes a significant host galaxy flux contribution from which we derive a spectroscopic redshift. The photometry of the host is suggestive of a high-redshift star-forming galaxy such as an LBG. The accepted method of







**Figure 1.** HSC-grizy light curves of HSC16adga (top), HSC17auzg (middle), and HSC17dbpf (bottom). Triangles indicate upper limits and errors are  $1\sigma$ . The rest-frame timescales are relative to the dates of detection, and the effective rest-frame wavelengths of the filters are given in the legends using  $z=2.399,\,1.965,\,$  and 1.851, respectively. Dates that spectra were acquired are identified by dashed black lines. *K*-corrected absolute magnitudes are discussed in M19.

<sup>&</sup>lt;sup>a</sup> COSMOS2015 unless otherwise specified.

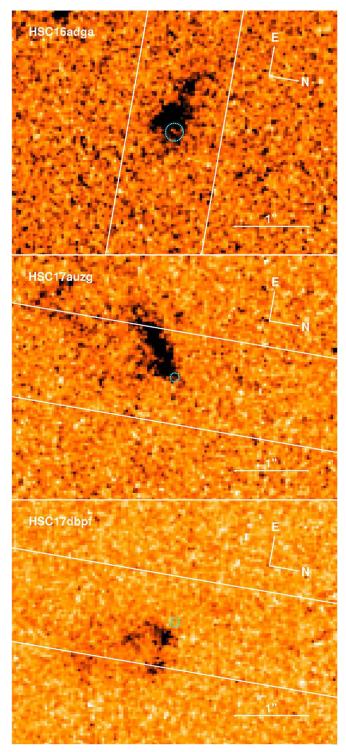
<sup>&</sup>lt;sup>b</sup> Rest frame, derived from Spec-z.

<sup>&</sup>lt;sup>c</sup> Derived using estimated wavelengths of emission features in raw 2D spectra.

 $<sup>^{</sup>m d}$  Two alternative MIZUKI photo-z peaks. COSMOS2015 photo-z not available.

<sup>&</sup>lt;sup>e</sup> No signal detected in the reduced imagery.

<sup>&</sup>lt;sup>16</sup> IRAF is distributed by the National Optical Astronomy Observatories, which are operated by the Association of Universities for Research in Astronomy, Inc., under cooperative agreement with the National Science Foundation.



**Figure 2.** Maps of the spectral slits used for the observations of HSC16adga (top), HSC17auzg (middle), and HSC17dbpf (bottom). The slits are projected onto publicly available *HST* images of the hosts in quiescence. The open circles mark the  $1\sigma$  regions of the transient locations as reported by SHIZUCA (see M19).

spectroscopic redshift determination of high-redshift LBGs is to infer an initial coarse redshift from visual confirmation of Ly $\alpha$ , the Ly $\alpha$  forest decrement, and strong interstellar medium (ISM) absorption lines, and to then constrain the redshift by comparison to LBG composite spectra with full FUV line lists and consideration of the behavior of the lines as a result of outflows (i.e., the ubiquitous presence of  $\sim 100-200 \, \mathrm{km \, s^{-1}}$ 

blueshifted ISM absorption line profiles and the P-Cygni-like profile of the Ly $\alpha$  line; Steidel et al. 1996, 1998, 2003, 2004; Pettini et al. 2000, 2001; Adelberger et al. 2003; Shapley et al. 2003; Cooke et al. 2006; Jones et al. 2012).

In the reduced spectrum of HSC16adga, we identify a strong absorption feature at  ${\sim}4100\,\mathring{\rm A},$  observer frame, with a slight flux break in the continuum at shorter wavelengths, which we attribute to  ${\rm Ly}\alpha$  absorption and the  ${\rm Ly}\alpha$  forest, respectively. For comparison, throughout we use the LBG composite spectra of Shapley et al. (2003). From Figure 3 several features are apparent (e.g., Si II  $\lambda1260$ , O I  $\lambda1302$ , C II  $\lambda1335$ , Si IV  $\lambda1394$ , O IV]  $\lambda1403$ ), indicating a redshift of  $z=2.399\pm0.004$ , consistent with both the COSMOS2015 photometric redshift estimate and the MIZUKI redshift probability distribution. The error is given in terms of the spectral resolution of LRIS, doubled to account for a 2 pixel smoothing algorithm applied to clarify absorption features.

At present, spectroscopic cross-correlation is ineffective for the positive determination of precise redshifts of high-redshift LBGs. One reason for this is that none of the identifiable FUV features of LBGs are at rest with respect to the redshift of the LBG system itself (Steidel et al. 2003). Still, cross-correlation can be used to quickly rule out the possibility of an alternative low-redshift fit.

We perform a simple cross-correlation of the spectrum with the Manual and Automatic Redshifting Software (MARZ; Hinton et al. 2016). MARZ cross-correlates input spectra with a library of galaxy templates and outputs the five most likely fits. Each fit includes a likelihood score from 1 to 4 representing respectively a likelihood of less than 50%, 50%–90%, 90%–99%, and greater than 99%. The MARZ top five redshift fits to the spectrum of HSC16adga are in the range 0.2845 < z < 2.42 with each fit receiving a score of 1, less than 50% likely.

As an independent check for the visual redshift estimates, we perform a coarse fit on the observed spectra with a library of spectral templates to infer their redshifts. The strong night sky lines are not always subtracted perfectly in spectroscopic observations, and the sky residuals are the main source of systematic uncertainties in observed spectra, which are not captured by noise spectra. To suppress such systematics, we bin the spectra in 200 Å windows (although in practice the inferred probability distributions are not very sensitive to the choice of bin size) by first clipping the top and bottom 15% of the flux distribution in each bin and then taking the weighted average of the remaining data points. The clipping here is intended to largely eliminate the systematics. We then feed the binned fluxes to the photo-z code, MIZUKI, to infer narrowband redshift probability distributions to complement the photometric redshifts derived from broadband photometry (see M19). Only galaxy templates are used in the fitting, and the supernova contribution to the spectrum is a source of uncertainty.

The MIZUKI narrowband redshift estimate of the spectrum of HSC16adga is shown in Figure 4. Compared to the broadband estimate, the probability of the spectroscopic redshift is higher because several alternative high-redshift probabilities have been eliminated (see M19).

# 3.2. HSC17auzg

HSC17auzg (SN-2016jhn) was first detected by SHIZUCA on MJD 57745 at coordinates (R.A., decl.) = (09:59:00.42, +02:14:20.8). The transient is detected in a star-forming host

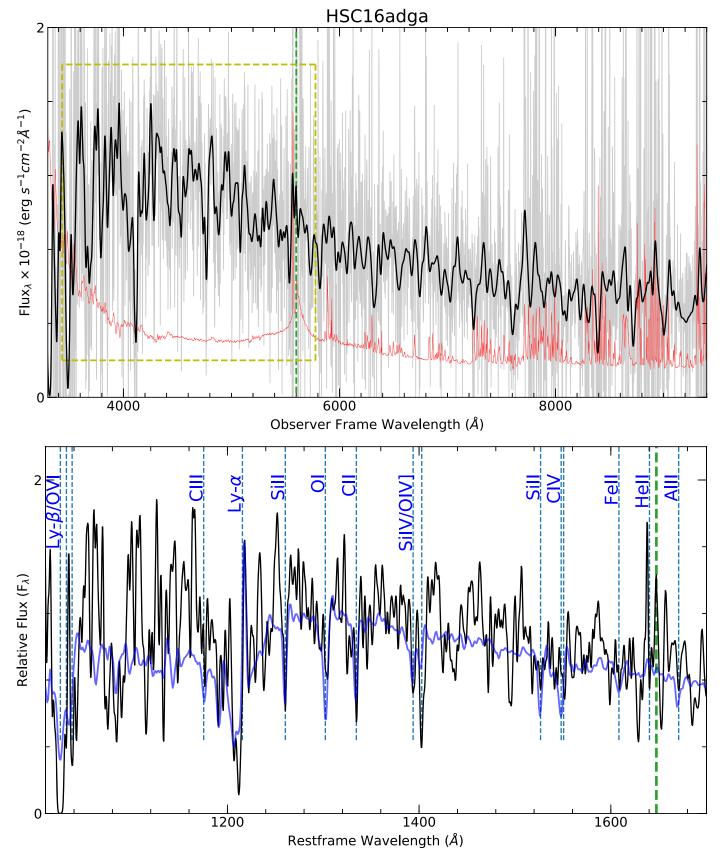
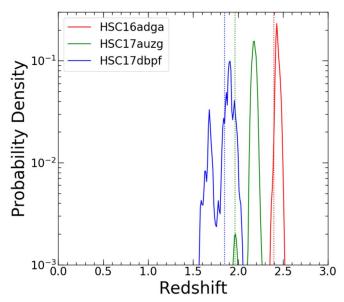


Figure 3. Top: flux-calibrated observer-frame 1D spectrum of HSC16adga, shown in gray. A smoothed spectrum is overlaid in black. The  $1\sigma$  pix<sup>-1</sup> error is shown in red. The green dashed line marks the dichroic separation at 5600 Å. The shape of the spectrum becomes less reliable shortward of 3600 Å (1060 Å rest frame) due to systematics in flat-fielding. Longward of 7000 Å (2060 Å rest frame), the spectrum becomes more sky-dominated and the smoothed spectrum has been manually clipped. The yellow dashed box is the zoomed-in region shown in the lower plot. Bottom: the same spectrum zoomed-in, smoothed, and redshifted to z = 2.399. A weak Ly $\alpha$  absorber LBG composite spectrum (see text) is overlaid in blue and scaled arbitrarily to emphasize the alignment of features. A subset of the features used to constrain the redshift is illustrated with dashed blue lines and labeled.



**Figure 4.** MIZUKI narrowband redshift probability distributions for each event. The MIZUKI broadband photometric redshift probability distributions are presented in identical fashion in M19.

with a COSMOS photo-z of  $1.65^{+0.06}_{-0.08}$  and a MIZUKI photo-z peaking at 1.78. It is observed offset by  $0.78 \pm 0.051$  ( $\sim 6$  kpc) from the host galaxy flux centroid, helping to rule out an AGN or a TDE. The light curve of HSC17auzg is shown in Figure 1.

We collected the Keck LRIS spectrum of HSC17auzg on MJD 57835.80. The positioning of the spectral slit is mapped in Figure 2. From the light curve in Figure 1 we estimate that the spectrum was taken some days after the rest-frame UV peak. The spectrum, composed of seven exposures (8400 s blue, 8253 s red), was reduced using standard IRAF procedures. The 1D spectrum is shown in Figure 5.

We identify a strong absorption feature at  $\sim 3500\,\text{Å}$ , observer frame, which we attribute to Ly $\alpha$ . From this we make a first approximation of the redshift whereby several key ISM absorption features of LBGs become apparent (e.g., O I  $\lambda$ 1302, Si IV  $\lambda$ 1394, O IV]  $\lambda$ 1403, Si II  $\lambda$ 1527, C IV  $\lambda\lambda$ 1548, 1551, Al II  $\lambda$ 1671), indicating a redshift of  $z=1.965\pm0.004$  (see Figure 5). This is significantly higher than the COSMOS2015 photometric redshift estimate but is represented in the MIZUKI redshift probability distribution.

The MARZ five most likely cross-correlation redshift fits to the spectrum of HSC17auzg are in the range 0.428 < z < 3.95 with each fit receiving a score of 1, less than 50% likely.

We produce a MIZUKI narrowband redshift estimate for HSC17auzg (see Figure 4). Compared to the broadband estimate (see M19), the probability of the spectroscopic redshift is similar, but its proximity to a probability peak is significantly improved. There is a more dominant probability peak, but spectroscopic feature alignment occurs only at the redshift of the secondary probability peak.

#### 3.3. HSC17dbpf

HSC17dbpf (SN-2017fei) was first detected by SHIZUCA on MJD 57816 at coordinates (R.A., decl.) = (09:58:33.42, +01:59:29.7). The transient is detected in a star-forming host with a COSMOS photo-z of  $2.25^{+0.08}_{-0.53}$  and a MIZUKI photo-z peaking at 1.58. It is observed offset by 0."58  $\pm$  0."052

( $\sim$ 5 kpc) from the host galaxy flux centroid, helping to rule out an AGN or a TDE. The light curve of HSC17dbpf is shown in Figure 1.

We collected two Keck LRIS spectra of HSC17dbpf on MJD 57835.93 and MJD 57836.84. The positioning of the spectral slit is mapped in Figure 2. The light curve in Figure 1 indicates that the event evolved very quickly, making the timing of the rest-frame UV peak difficult to estimate, but it appears that the spectra were taken within a few days of this peak. The spectra (see Figure 6) were collected over two consecutive nights, which we treat as a single epoch composed of five exposures (6000 s blue, 5895 s red).

We attribute the strong absorption feature at  $\sim 3400\,\text{Å}$ , observer frame, to Ly $\alpha$ . From this we make a first approximation of the redshift whereby several key ISM absorption features of LBGs become apparent (e.g., C II  $\lambda 1335$ , Si IV  $\lambda 1394$ , O IV]  $\lambda 1403$ , Si II  $\lambda 1527$ , C IV  $\lambda \lambda 1548$ , 1551), indicating a redshift of  $z=1.851\pm0.004$  (see Figure 6). This is consistent with the COSMOS2015 photometric redshift estimate but higher than predicted by the MIZUKI redshift probability distribution.

The MARZ five most likely cross-correlation redshift fits to the spectrum of HSC17dbpf are in the range 0.0155 < z < 4.67 with each fit receiving a score of 1, less than 50% likely.

We produce a MIZUKI narrowband redshift estimate for HSC17dbpf (see Figure 4). Compared to the broadband estimate, which is inconsistent with the spectroscopic redshift (see M19), the probable redshift range is similar but encompasses higher values including the spectroscopic redshift.

## 4. Discussion

## 4.1. Nature

The high-redshift supernova nature of each event is well supported. The spectroscopic redshifts are precisely constrained and consistent with the MIZUKI narrowband redshifts, which offer no low-redshift probability, and cross-correlation with MARZ produces no reliable low-redshift alternative. Furthermore, when the measured redshifts are applied to the transient photometry, M19 find low-redshift supernovae with comparable UV light curves in each case.

Although TDEs can resemble supernovae in their light curve behavior, these are highly centralized events, and our candidates are all significantly offset from their host-galaxy flux centroids (see Figure 2). Thus it is unlikely that the candidates are of this type. In Figure 7 we compare the observed spectra to the FUV spectrum of the TDE ASASSN-14li (Cenko et al. 2016). We include a Sloan Digital Sky Survey (SDSS) quasi-stellar object (QSO, equivalent to an AGN) composite spectrum (Vanden Berk et al. 2001) to which the TDE is similar. The lack of any strong N V emission or any other nitrogen emission in the observed spectra is a notable difference to that expected of a TDE.

The SLSN class of supernovae was first defined as a luminosity class that includes all supernovae exceeding an absolute magnitude of M=-21 in any band (Gal-Yam 2012). However, currently SLSNe are identified and grouped spectroscopically and comprise events with a continuum of peak magnitudes that extends below the M=-21 threshold (Nicholl et al. 2017; Moriya et al. 2018; Quimby et al. 2018).

The k-corrected peak absolute magnitudes of the presented SHIZUCA photometric candidates (see M19) approach but do

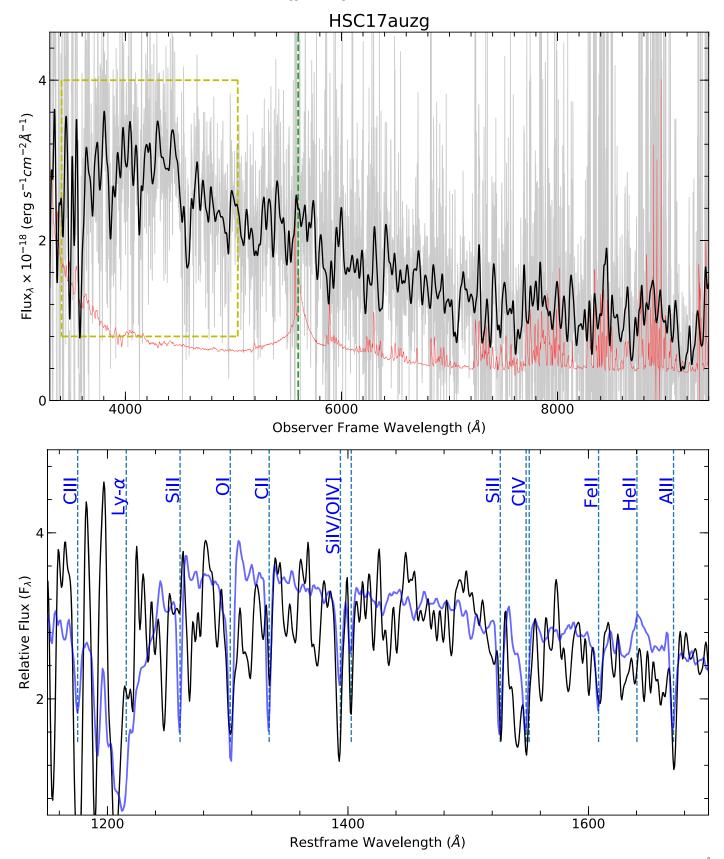


Figure 5. Top: flux-calibrated observer-frame 1D spectrum of HSC17auzg, similar to Figure 3. The shape of the spectrum becomes less reliable shortward of 3600 Å (1215 Å rest frame), due to systematics in flat-fielding. Longward of 7000 Å (2360 Å rest frame), the spectrum becomes more sky-dominated. Bottom: the same spectrum zoomed-in, smoothed, and redshifted to z = 1.965, similar to Figure 3. A strong Ly $\alpha$  absorber LBG composite spectrum (see text) is overlaid in blue and scaled arbitrarily to emphasize the alignment of features. A subset of the features used to constrain the redshift is illustrated.

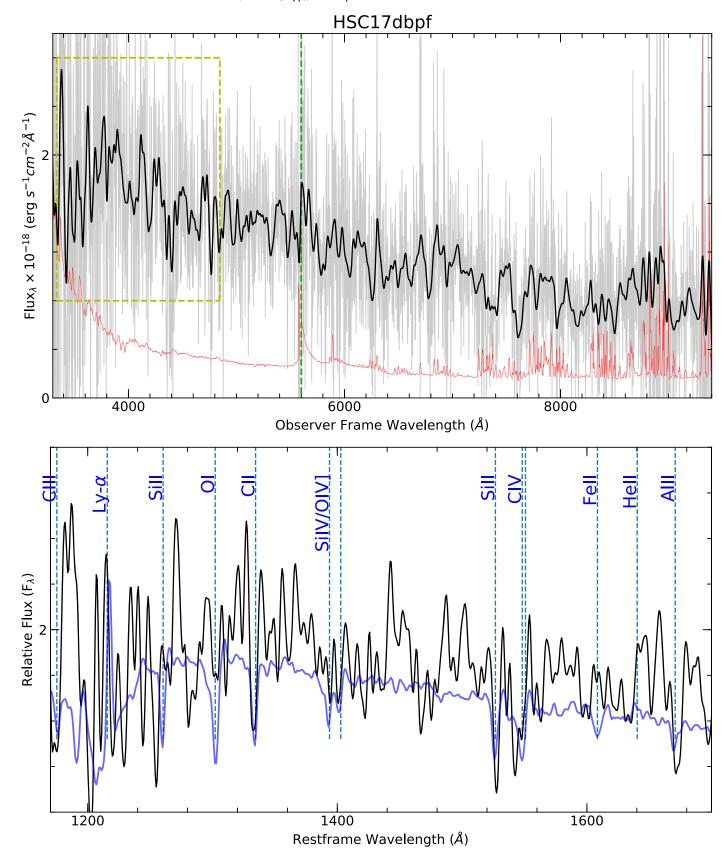


Figure 6. Top: flux-calibrated observer-frame 1D spectrum of HSC17dbpf, similar to Figure 3. The shape of the spectrum becomes less reliable shortward of 3600 Å (1260 Å rest frame) due to systematics in flat-fielding. Longward of 7000 Å (2455 Å rest frame), the spectrum become more sky-dominated. Bottom: the same spectrum zoomed-in, smoothed, and redshifted to z=1.851, similar to Figure 3. A weak Ly $\alpha$  absorber LBG composite spectrum (see text) is overlaid in blue and scaled arbitrarily to emphasize the alignment of features. A subset of the features used to constrain the redshift is illustrated.

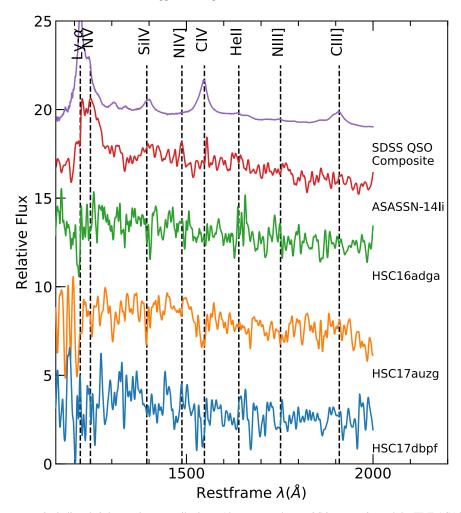


Figure 7. All three supernova spectra including their host galaxy contributions. Also presented are a QSO composite and the TDE ASASSN-14li. We have increased the noise in the spectrum of the TDE to approximate that of our observations, and we have applied a Ly $\alpha$  deforestation factor to the TDE and the QSO at wavelengths shortward of 1216 Å consistent with a redshift of  $z \sim 2$ . The labeled transitions are emission features observed in ASASSN-14li by Cenko et al. (2016), most of which are also commonly seen in AGNs (as is evident from the QSO composite).

not meet the strict photometric criteria of SLSNe set by Gal-Yam (2012). However, the overluminous aspect of each of the candidates in the FUV, while difficult to reconcile with ordinary supernovae, is representative of SLSNe (Quimby et al. 2011; Yan et al. 2017b, 2018). The intermediate FWHM seeing and light cirrus during spectroscopic follow-up resulted in the candidate spectra having lower S/N than anticipated and prevents positive supernova type determinations. Thus, as implied throughout, we identify these transients as probable SLSNe. Late-time follow-up of the candidates could enable supernova type determinations by providing more precise host-galaxy subtraction templates and revealing any late-time circumstellar emission features as seen in some SLSNe (Fox et al. 2015; Yan et al. 2017a).

#### 4.2. Host Subtraction

We compare our host-subtracted spectra to some of the few FUV spectra of early-time SLSNe and UV-bright SNe in an effort to discern the spectroscopic type of each supernova. While unsuccessful in this regard, the comparisons still offer useful information.

Because we do not have spectra of the host galaxies alone, we perform a "pseudo-host" subtraction on each spectrum to

enable our analysis. We assume an LBG-like host for each supernova, consistent with the templates fit with MIZUKI to the broadband photometry (see M19) and with the hosts of SLSNe in general. Analyses of a large sample of LBGs (Shapley et al. 2003; Cooke et al. 2009) have shown that the characteristics and colors of an LBG spectrum are strongly correlated with the equivalent width of the Ly $\alpha$  feature. Shapley et al. (2003) divided their LBG sample into quartiles based on this equivalent width and subsequently constructed composite spectra for each. The quartiles are referred to here as strong Ly $\alpha$  absorbers, weak Ly $\alpha$  absorbers, weak Ly $\alpha$  emitters, and strong Ly $\alpha$  emitters. We include these composites in our pseudo-host constructions using the estimated Ly $\alpha$  absorption or emission strengths of the hosts to select the specific composite for use in each subtraction.

The wavelength ranges of the spectra include the HSC-*gri* bands. To construct the pseudo-hosts over this whole range, we use the LBG composite spectra from 912 to 2000 Å (their full extent) and starburst (SB) templates from Calzetti et al. (1994) at wavelengths beyond 2000 Å. This extension is valid because LBGs are highly star-forming galaxies. Ly $\alpha$  absorbers have redder UV continua than Ly $\alpha$  emitters (Shapley et al. 2003; Cooke et al. 2009). We therefore use the SB2 template (0.10 < E(B-V) < 0.21) to complete a pseudo-host constructed with a

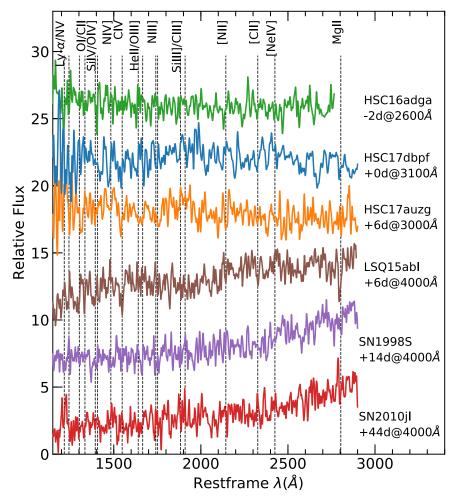


Figure 8. Three pseudo-host-subtracted supernova spectra (see text) along with the SLSN-II LSQ15abl and two UV-bright SNe-IIn. The spectra are labeled along with their approximate phase in days since peak at the associated rest-frame wavelength. We have increased the noise in the comparison spectra to approximate that of our observations, and we have applied a Ly $\alpha$  forest effect for a redshift of  $z\sim2$  at wavelengths shortward of 1216 Å. The labeled transitions are emission features identified in SNe-IIn by Fransson et al. (2005, 2014).

Ly $\alpha$  absorber LBG composite, and we use the SB1 template (E(B-V) < 0.10) to complete a pseudo-host constructed with a Ly $\alpha$  emitter LBG composite.

To accurately scale the pseudo-hosts, we calculate the ratio of host flux to supernova flux in each spectrum using the associated HSC photometry as measured most closely in time to our observations (see M19). The high-redshift hosts are nearpoint sources, and the bulk of the flux enters the spectroscopic slit in each case. But each slit is centered on the supernova, so some of the host flux may not be included (see Figure 2). When necessary, we apply a geometrical correction to account for any portion of the host not within the spectroscopic slit.

## 4.3. SLSNe-II

Type-II supernovae (SNe-II) are defined as those that exhibit hydrogen lines in spectra, typically H- $\alpha$  (Filippenko 1997), though they can also be confirmed with Ly $\alpha$  emission in FUV spectra (Fransson et al. 2002, 2005, 2014). LBGs, which are highly star-forming galaxies, are a common host type for high-redshift supernovae (Cooke et al. 2009, 2012). But they are also often strong Ly $\alpha$  absorbers or emitters, which complicates the task of identifying supernova Ly $\alpha$  emission.

Observed Ly $\alpha$  emission in a high-redshift supernova/host spectrum can still be used as a positive identifier of a SNe-II if

it can be confidently attributed to the supernova. In cases where the host and supernova are spatially separated on the sky such that their profiles are distinguishable, the source or sources of Ly $\alpha$  emission can be observed directly. If this is not the case, the source of observed Ly $\alpha$  emission can still be inferred from the characteristics of the feature. Ly $\alpha$  emission from a supernova originates in the ejecta or a circumstellar medium (CSM) or both. Ly $\alpha$  emission from ejecta is broadened according to the ejecta velocity. Ly $\alpha$  from a cold, slowmoving CSM is observed as a narrow feature with a blue velocity offset from the redshift of the host. In cases of LBGs with Ly $\alpha$  in emission, the feature arises as Ly $\alpha$  emission from the LBG is back-scattered off receding gas outflows and passes back through the LBG off-resonance avoiding absorption. Consequently, the (narrow) feature is typically redshifted from the host by up to  $1000 \,\mathrm{km \, s^{-1}}$  or more (Shapley et al. 2003), thus distinguishing it from any relatively blueshifted supernova Ly $\alpha$  emission.

High-redshift supernova  $Ly\alpha$  emission can also be largely or completely absorbed by neutral hydrogen (H I) in the host interstellar medium, its circumgalactic medium, and the intervening intergalactic medium. So the absence of observable supernova  $Ly\alpha$  emission does not necessarily imply the absence of hydrogen integral to the explosion scenario. In

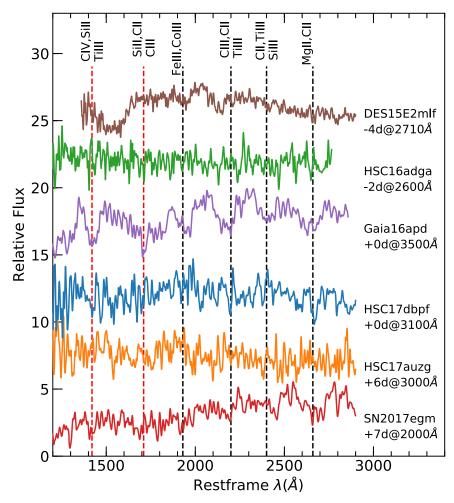


Figure 9. Three pseudo-host-subtracted supernova spectra (see text) along with the high-redshift SLSN-I, DES15E2mlf, and the two low-redshift SLSNe-I. The spectra are labeled along with their approximate phase in days since peak at the associated rest-frame wavelength. We have increased the noise of the comparison spectra to approximate that of our observations, and have applied a Ly $\alpha$  forest effect for a redshift of  $z \sim 2$  at wavelengths shortward of 1216 Å. Transitions reproduced in the spectral models of Mazzali et al. (2016) are marked with black dotted lines, while those marked with red dotted lines are derived by Yan et al. (2018) with a model fitting method.

cases where the lack of Ly $\alpha$  emission is ambiguous, the presence of certain features unattributable to the host alone may be another reliable way to identify an SLSN-II from its FUV spectrum. For example, most SLSNe-II are spectroscopically classified as SNe-IIn based on the presence of narrow hydrogen lines in emission (Ofek et al. 2007; Smith et al. 2007, 2008; Drake et al. 2010; Chatzopoulos et al. 2011). Studies by Fransson et al. (2005, 2014) have shown that SN-IIn FUV spectra can have signature strong emission features such as O I  $\lambda$ 1300, N IV  $\lambda$ 1486, and Mg II  $\lambda$ 2800. These features could conceivably be used to confirm an SLSN-II classification (Cooke et al. 2009). Magnesium emission would be a particularly useful identifying feature of SLSNe-II as it has been observed to persist strongly to very late times in SNe-IIn (Fransson et al. 2005).

In Figure 8 we compare our pseudo-host-subtracted spectra to the early-time FUV spectrum of the low-redshift SLSN-II LSQ15abl (Quimby et al. 2015) and the earliest FUV spectra of the UV-luminous SNe-IIn SN1998S (Fransson et al. 2005) and SN2010jl (Fransson et al. 2014). None of our supernovae exhibit obvious Ly $\alpha$  emission. It is tempting to attribute the blunted emission-like feature at  $\sim$ 1190 Å in the spectrum of HSC17dbpf to a component of slightly blueshifted Ly $\alpha$ 

emission, but this feature arises at a wavelength at which the shape of the spectrum is unreliable due to systematics in flat-fielding. It is worth noting that in the noise-added spectra being used for comparison,  $Ly\alpha$  emission is likewise not obvious. We know that these events do exhibit  $Ly\alpha$  emission, so this comparison illustrates the baseline S/N required of a high-redshift supernova spectrum in order to rule out at least any unmolested  $Ly\alpha$  emission. The coarse blackbody shapes of our spectra are bluer than the comparison spectra, but this is not surprising given their earlier collection times. Beyond these observations, the spectra are simply of too low S/N to say anything about the more subtle emission features anticipated in the FUV of SLSNe-II.

#### 4.4. SLSNe-I

The spectra of SLSNe-I are uniquely conspicuous, displaying a number of characteristic broad hump-like features throughout the optical and UV (Quimby et al. 2011; Howell et al. 2013; Nicholl et al. 2015; Yan et al. 2017b, 2018). The number of spectroscopically observed SLSNe-I is increasing rapidly, and the nature of these features is being clarified (Mazzali et al. 2016; Lunnan et al. 2018; Quimby et al. 2018).

Our understanding of the UV features of SLSNe-I is also improving thanks in large part to spectroscopic studies of highredshift examples (Berger et al. 2012; Howell et al. 2013; Pan et al. 2017; Smith et al. 2018). From what has been observed to date, SLSNe-I typically exhibit broad absorption troughs centered around  $\sim$ 1930, 2200, 2400, and 2660 Å (Smith et al. 2018). Spectral synthesis models have attributed these features to blended transitions of iron, cobalt, carbon, titanium, silicon, and magnesium (Mazzali et al. 2016). Prominent features at still shorter wavelengths (e.g., 1400, 1700 Å) have been observed in spectra of high-redshift SLSNe-I for which UV coverage extends sufficiently far (Pan et al. 2017), as well as in low-redshift SLSNe-I with HST FUV spectral coverage (Yan et al. 2017b, 2018). But the consistency of these features among such events is insufficient to draw conclusions about their normalcy at this time.

In Figure 9 we compare our pseudo-host-subtracted spectra to an extensive FUV spectrum of the high-redshift SLSN-I DES15E2mlf (Pan et al. 2017) and the two low-redshift SLSNe-I with FUV spectral coverage, Gaia16apd (Yan et al. 2017b) and SN2017egm (Yan et al. 2018). In the noise-added comparison spectra, the absorption troughs are significantly dampened, but most are still discernible. Differences in the strengths, widths, and central wavelengths of individual troughs are also apparent. The longer wavelength features are not observed to any significance in our high-redshift sample, though at these wavelengths our spectra are all sky-dominated. There seem to be some similarities between the comparison spectra and the spectra of HSC17auzg and HSC17dbpf at short wavelengths (e.g., the absorption dip at ~1400 Å, the general trends of the spectra from  $\sim$ 1700 to 2000 Å). Positive SLSN-I classifications can be made with optical spectra by performing a  $\chi^2$  minimization of model spectra created by scaling, adding, and then redshifting supernova and host galaxy templates (Howell et al. 2005; Quimby et al. 2018). We attempt to quantify the significance of any similarities between our spectra and the FUV spectra of SLSNe-I by the same method. However, with so few FUV spectra of SLSNe-I available for comparison, the significance of the presence or lack of features in the FUV spectra of our targets cannot be quantified, and our results are inconclusive.

### 5. Conclusion

With this SHIZUCA pilot spectroscopic program, we have demonstrated the feasibility of spectroscopic follow-up of high-redshift SLSNe, testing the current limits of high-redshift transient spectral analysis. The high éntendu and dense cadence of the HSC-SSP combined with the detection efficiency of SHIZUCA can generate high-redshift SLSN candidates at a sufficient rate to statistically enable follow-up of several targets near maximum light at any time during an active observation campaign. Keck LRIS provides the best facility for spectroscopic follow-up of  $z\sim 2$  supernovae in terms of light-gathering power combined with blue sensitivity.

We have reported on FUV spectroscopic follow-up of three probable SLSNe at  $z=1.851,\ 1.965,\$ and 2.399. The spectra are extracted from  $\sim 2$  hr exposures taken under suboptimal seeing and weather conditions. The S/N of the reduced spectra are relatively low, and flux measurements become unreliable at wavelengths shortward of 3600 Å, observer frame. Thus in LRIS spectra of similar S/N, moderate-to-weak Ly $\alpha$  features, perhaps dampened by H I absorption, cannot be confirmed or

ruled out at  $z \lesssim 2$ . Such analyses can be extended to shorter wavelengths and lower redshifts using spectra with higher S/N or exhibiting strong Ly $\alpha$  emission.

Our analysis of the pseudo-host-subtracted supernova spectra is inconclusive regarding the spectroscopic type of each event. We estimate that doubling the S/N of the spectra presented here would enable supernova type confirmations of future events by substantiating any strong emission features of SLSNe-II or the broad absorption troughs of SLSNe-I. By observing under optimal seeing conditions, employing longer exposures, and perhaps taking advantage of brighter than average targets, this S/N is achievable.

Future follow-up is merited to ascertain if any of the supernovae exhibit late-time strong emission (e.g., Ly $\alpha$ , Mg II) from which type information can be inferred. In addition, late-time spectra of the hosts with little-to-no supernova contributions can be used to produce more precise host-subtracted supernova spectra and better constrain their UV continuum slopes, enabling further analysis.

An approximate rate calculation of  $z\sim 2$  SLSNe based on the SHIZUCA photometric candidate catalog is discussed in M19. The accuracy of this rate is directly dependent on the efficiency of SHIZUCA in selecting high-redshift transients and reducing photometric redshift confusion. More spectra are needed to establish and improve this efficiency and separate high-redshift supernovae from other high-redshift transient phenomena such as AGN and TDEs. As a first approximation, three of five SHIZUCA photometric candidates with spectroscopically confirmed redshifts (excluding the nondetected HSC17davs) are at high redshift, which gives a current redshift efficiency of 60%.

SHIZUCA is capable of detecting supernovae to z > 6, and with the right observing strategy and clear skies, spectroscopic follow-up of these events out to the edge of the epoch of reionization may be achievable. The potential of the *James Webb Space Telescope* will be greater still, capable of acquiring spectra of SLSNe as far as z = 20. At such high redshifts, only deep, wide-field infrared surveys will be capable of producing targets. Such surveys are already being considered in future facilities like the University of Tokyo Atacama Observatory (TAO) and the Kunlun Dark Universe Survey Telescope (KDUST). By exploring the current limits of high-redshift transient astronomy, SHIZUCA and similar programs are setting the stage for observing the explosions of the very first stars.

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References
Adelberger, K. L., Steidel, C. C., Shapley, A. E., & Pettini, M. 2003, ApJ,
Aihara, H., Armstrong, R., Bickerton, S., et al. 2018, PASJ, 70, 8
Angus, C. R., Levan, A. J., Perley, D. A., et al. 2016, MNRAS, 458, 84
Bellm, E., & Kulkarni, S. 2017, NatAs, 1, 0071
Berger, E., Chornock, R., Lunnan, R., et al. 2012, ApJL, 755, L29
Bufano, F., Immler, S., Turatto, M., et al. 2009, ApJ, 700, 1456
Calzetti, D., Kinney, A. L., & Storchi-Bergmann, T. 1994, ApJ, 429, 582
Cenko, S. B., Cucchiara, A., Roth, N., et al. 2016, ApJL, 818, L32
Chatzopoulos, E., Wheeler, J. C., Vinko, J., et al. 2011, ApJ, 729, 143
Cooke, J., Sullivan, M., Barton, E. J., et al. 2009, Natur, 460, 237
Cooke, J., Sullivan, M., Gal-Yam, A., et al. 2012, Natur, 491, 228
Cooke, J., Wolfe, A. M., Gawiser, E., & Prochaska, J. X. 2006, ApJ, 652, 994
Dark Energy Survey Collaboration, Abbott, T., Abdalla, F. B., et al. 2016,
            460, 1270
Drake, A. J., Djorgovski, S. G., Prieto, J. L., et al. 2010, ApJL, 718, L127
Filippenko, A. V. 1997, ARA&A, 35, 309
Fox, O. D., Smith, N., Ammons, S. M., et al. 2015, MNRAS, 454, 4366
Fransson, C., Challis, P. M., Chevalier, R. A., et al. 2005, ApJ, 622, 991
Fransson, C., Chevalier, R. A., Filippenko, A. V., et al. 2002, ApJ, 572, 350
Fransson, C., Ergon, M., Challis, P. J., et al. 2014, ApJ, 797, 118
Gal-Yam, A. 2012, Sci, 337, 927
Hinton, S. R., Davis, T. M., Lidman, C., Glazebrook, K., & Lewis, G. F. 2016,
Howell, D. A., Kasen, D., Lidman, C., et al. 2013, ApJ, 779, 98
Howell, D. A., Sullivan, M., Perrett, K., et al. 2005, ApJ, 634, 1190
Jones, T., Stark, D. P., & Ellis, R. S. 2012, ApJ, 751, 51
Kaiser, N., Burgett, W., Chambers, K., et al. 2010, Prov. SPIE, 7733, 77330E
Laigle, C., McCracken, H. J., Ilbert, O., et al. 2016, ApJS, 224, 24
Leloudas, G., Schulze, S., Krühler, T., et al. 2015, MNRAS, 449, 917
Lunnan, R., Chornock, R., Berger, E., et al. 2014, ApJ, 787, 138
Lunnan, R., Chornock, R., Berger, E., et al. 2018, ApJ, 852, 81
Madau, P., & Dickinson, M. 2014, ARA&A, 52, 415
Mazzali, P. A., Sullivan, M., Pian, E., Greiner, J., & Kann, D. A. 2016,
            , 458, 3455
Moriya, T. J., Sorokina, E. I., & Chevalier, R. A. 2018, SSRv, 214, 59
Moriya, T. J., Tanaka, M., Yasuda, N., et al. 2019, ApJS, 241, 16
Mould, J., Abbott, T., Cooke, J., et al. 2017, Sci. Bull., 62, 675
Nicholl, M., Guillochon, J., & Berger, E. 2017, ApJ, 850, 55
Nicholl, M., Smartt, S. J., Jerkstrand, A., et al. 2015, MNRAS, 452, 3869
Ofek, E. O., Cameron, P. B., Kasliwal, M. M., et al. 2007, ApJL, 659, L13
```

Oke, J. B., Cohen, J. G., Carr, M., et al. 1995, PASP, 107, 375

Pan, Y.-C., Foley, R. J., Smith, M., et al. 2017, MNRAS, 470, 4241

```
Panagia, N. 2007, in AIP Conf. Ser. 937, Supernova 1987A: 20 Years After:
   Supernovae and Gamma-Ray Bursters, ed. S. Immler, K. Weiler, &
   R. McCray (Melville, NY: AIP), 236
Pettini, M., Shapley, A. E., Steidel, C. C., et al. 2001, ApJ, 554, 981
Pettini, M., Steidel, C. C., Adelberger, K. L., Dickinson, M., & Giavalisco, M.
   2000, ApJ, 528, 96
Prajs, S., Sullivan, M., Smith, M., et al. 2017, MNRAS, 464, 3568
Quimby, R. 2014, Cycle 22, HST Proposal 13784 (Baltimore, MD: STSci)
Quimby, R. 2016, Cycle 24, HST Proposal 14674 (Baltimore, MD: STSci)
Quimby, R. M., Cooke, J., & Pritchard, T. 2015, ATel, 7439
Quimby, R. M., De Cia, A., Gal-Yam, A., et al. 2018, ApJ, 855, 2
Quimby, R. M., Kulkarni, S. R., Kasliwal, M. M., et al. 2011, Natur, 474, 487
Quimby, R. M., Yuan, F., Akerlof, C., & Wheeler, J. C. 2013, MNRAS,
   431, 912
Schulze, S., Krühler, T., Leloudas, G., et al. 2018, MNRAS, 473, 1258
Scoville, N., Aussel, H., Benson, A., et al. 2007, ApJS, 172, 150
Shapley, A. E., Steidel, C. C., Pettini, M., & Adelberger, K. L. 2003, ApJ,
```

588, 65

```
Shull, J. M. 2009, in AIP Conf. Ser. 1135, Future Directions in Ultraviolet
  Spectroscopy, ed. M. E. van Steenberg (Melville, NY: AIP), 301
Smith, M., Sullivan, M., D'Andrea, C. B., et al. 2016, ApJL, 818, L8
Smith, M., Sullivan, M., Nichol, R. C., et al. 2018, ApJ, 854, 37
Smith, N., Chornock, R., Li, W., et al. 2008, ApJ, 686, 467
Smith, N., Li, W., Foley, R. J., et al. 2007, ApJ, 666, 1116
Steidel, C. C., Adelberger, K. L., Dickinson, M., Giavalisco, M., & Pettini, M.
   1998, ApJ, 492, 428
Steidel, C. C., Adelberger, K. L., Shapley, A. E., et al. 2003, ApJ, 592, 728
Steidel, C. C., Giavalisco, M., Pettini, M., Dickinson, M., & Adelberger, K. L.
  1996, ApJL, 462, L17
Steidel, C. C., Shapley, A. E., Pettini, M., et al. 2004, ApJ, 604, 534
Tanaka, M. 2015, ApJ, 801, 20
Vanden Berk, D. E., Richards, G. T., Bauer, A., et al. 2001, AJ, 122, 549
Woodgate, B. E., Kimble, R. A., Bowers, C. W., et al. 1997, Prov. SPIE, 3118, 2
Yan, L., Lunnan, R., Perley, D. A., et al. 2017a, ApJ, 848, 6
Yan, L., Quimby, R., Gal-Yam, A., et al. 2017b, ApJ, 840, 57
```

Yan, Y., Pereley, D., De Cia, A., et al. 2018, ApJ, 858, 91